

Transcritical Bifurcations and Algebraic Aspects of Quadratic Multiparametric Families

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ABSTRACT. This article reveals an analysis of the quadratic systems that hold multiparametric families therefore, in the first instance the quadratic systems are identified and classified in order to facilitate their study and then the stability of the critical points in the finite plane, its bifurcations, stable manifold and lastly, the stability of the critical points in the infinite plane, afterwards the phase portraits resulting from the analysis of these families are graphed. To properly perform this study it was necessary to use some results of the non-linear systems theory, for this reason vital definitions and theorems were included because of their importance during the study of the multiparametric families. Algebraic aspects are also included.

Keywords and Phrases. *Quadratic Polynomial Systems, Critical Points, Bifurcations, Stable Manifold, Phase portraits of polynomial systems.*

1. Introduction

Systems of differential equations are known to express a number of mathematical, physical and engineering situations. In particular, this article is based about the study of all quadratic multiparametric subfamilies associated with the next family: Given the family with $a, b, c, m, k \in \mathbb{R}$.

$$(1.1) \quad \begin{cases} \dot{x} &= y \\ \dot{y} &= (\alpha x^{m+k-1} + \beta x^{m-k-1})y - \gamma x^{2m-2k-1} \end{cases}$$

We can find antecedents of the algebraic and qualitative studies of this family in [1, 2, 3, 4]. Another algebraic and dynamical studies can be found in [5, 6]. In general, we can see qualitative studies about planar systems in [7], [8] and [9], furthermore antecedents of applied bifurcations study in [10]. In the present work, we take Proposition 4.1, pag 12, in [2, 3] which the goal of analyze each quadratic subfamily equivalently to (1.1). Considering the constants a, b, c , and $s, p, r \in \mathbb{Z}^+$. Then, we analyze different cases to determine quadratic systems attached to (1.1) taking into account the regions in the space determined by the for the different parameters.

For the study of the quadratic multiparametric families described by (1.1) where we use different topics studied in [11],[12],[13] and [14]. Then, we find the critical

points associated with each quadratic family and analyzing their stability in both the finite and infinite planes, also we present a deeper study determined by regions to see the changes in stability of the critical points and from here analyze bifurcations presented in some families. Finally, we used software like [14] and [15] for the detailed construction of the behaviors by means of the global phase portrait associated with each quadratic multiparametric family.

2. Preliminaries

In this section we provide the necessary theoretical background to understand the rest of the paper.

A planar polynomial system of degree n is given by:

$$(2.1) \quad \begin{aligned} \dot{x} &= P(x,y) \\ \dot{y} &= Q(x,y) \end{aligned}$$

Where $P, Q \in \mathbb{C}[x, y]$, and n is given by $n = \max(\deg P, \deg Q)$

We denote the polynomial vector field associated to the system (2.1) like: By $X := (P, Q)$. The planar polynomial vector field X can be also written in the form:

$$X = P(x, y) \frac{\partial}{\partial x} + Q(x, y) \frac{\partial}{\partial y}.$$

A differential equations associate to polynomial vector field of the form (2.1) is given by:

$$\frac{dy}{dx} = \frac{Q(x,y)}{P(x,y)}.$$

THEOREM 2.1. *Hyperbolic Singular Points Theorem.*

Let $(0, 0)$ be an isolated singular point of the vector field X , given by,

$$(2.2) \quad \begin{cases} \dot{x} &= ax + by + A(x, y) \\ \dot{y} &= cx + dy + B(x, y) \end{cases}$$

Where A and B are analytic in a neighborhood of the origin with $A(0, 0) = B(0, 0) = DA(0, 0) = DB(0, 0) = 0$. Let λ_1 and λ_2 be the eigenvalues of the linear part $DX(0, 0)$ of the system at the origin. Then:

- (a) If λ_1, λ_2 are real and $\lambda_1 \lambda_2 < 0$, then $(0, 0)$ is a saddle, where separatrix call $(0, 0)$ in the directions given by the eigenvectors associated with λ_1 and λ_2 .
- (b) If λ_1, λ_2 are real and $\lambda_1 \lambda_2 > 0$, then $(0, 0)$ is a node. If $\lambda_1 > 0 (\lambda_1 < 0)$ then it is repelling or unstable (respectively attracting or stable).
- (c) If $\lambda_1 = \alpha + \beta_i$ y $\lambda_2 = \alpha - \beta_i$ with $\alpha, \beta \neq 0$ then $(0, 0)$ is a focus. If $\alpha > 0$ or $(\alpha < 0)$ it is repelling or unstable (respectively attracting or stable).
- (d) If $\lambda_1 = \beta_i$ and $\lambda_2 = -\beta_i$, then $(0, 0)$ is a linear center, focus or a center.

for a more detailed study, please see [13, pág 71]

THEOREM 2.2. Non-Hyperbolic Singular Points Theorem

Let $(0, 0)$ be an isolated singular point of the vector field X given by:

$$(2.3) \quad \begin{cases} \dot{x} = y + A(x, y) \\ \dot{y} = B(x, y) \end{cases}$$

Where X and Y are analytic in a neighborhood of the point $(0, 0)$ and considers the series expansion have expansions starting with terms of the second degree in x and y . Let $y = f(x) = a_2x^2 + a_3x^3 + \dots$ be the solution of the equation $y + A(x, y) = 0$ in a neighborhood of the point $(0, 0)$, and suppose you have the following series expansion of the function $F(x) = B(x, f(x)) = ax^m(1 + \dots)$ y $G(x) = (\frac{\partial A}{\partial x} + \frac{\partial B}{\partial y})(x, f(x)) = bx^n(1 + \dots)$ where $a \neq 0$, $m \geq 2$, and $n \geq 1$. Then:

- (1) If $G(x) \equiv 0$ and $F(x) = ax^m \dots$ for $m \in \mathbb{N}$ with $m \geq 1$ and $a \neq 0$, then:
 - (i) if m is odd and $a > 0$, then the origin of X is a saddle and If $a < 0$, then it is a center or a focus.
 - (ii) If m is even then the origin of X is a cusp.
- (2) If $F(x) = ax^m + \dots$ and $G(x) = bx^n + \dots$ with $m \in \mathbb{N}$, $m \geq 2$, $n \in \mathbb{N}$, $n \geq 1$, $a \neq 0$ and $b \neq 0$. Then we have:
 - (i) If m is even, and
 - (i.a) $m < 2n + 1$, then the origin of X is a cusp.
 - (i.b) $m > 2n + 1$, then the origin of X is a saddle-node.
 - (ii) If m is odd and $a > 0$, then the origin of X is a saddle.
 - (iii) If m is odd, $a < 0$ and
 - (iii.a) $m < 2n + 1$, or $m = 2n + 1$ and $b^2 + 4a(n + 1) < 0$, then the origin of X is a center or a focus.
 - (iii.b) n is odd and either $m > 2n + 1$, or $m = 2n + 1$ and $b^2 + 4a(n + 1) \geq 0$. Then the phase portrait of the origin of X consists of one hyperbolic and one elliptic sector.
 - (iii.c) n is even and either $m > 2n + 1$, or $m = 2n + 1$ and $b^2 + 4a(n + 1) \geq 0$. Then the origin of X is a node. The node is attracting if $b < 0$ and repelling if $b > 0$.

For a more detailed study, see [13, pág 116]

2.1. Bifurcations. We consider the system, depend on a parameter λ :

$$(2.4) \quad \dot{x} = f(x, \lambda)$$

If the change in λ that produces a qualitative or topological change in the behavior of the planar system (2.4), his is called a Bifurcations. This can be a local bifurcation occurs when the change in the parameter causes a change in the stability of an equilibrium point. Global bifurcations normally occur in larger invariant sets of the system.

Codimension - One Bifurcations These Bifurcations require the variation of a single parameter to occur in the system, all have a normal form, that is, a topologically equivalent system, either local or global to the initial system.

Transcritical Bifurcations: A transcritical Bifurcations in critical point exists for every value of the parameter λ but they exchange their stability with another critical point after the “collision” between them.

Saddle-focus-saddle Bifurcations:

DEFINITION 2.3. We Will call a Bifurcations saddle-focus-saddle is when a parameter change it implies that two critical points, one saddle, collapse in a focus and later they recover its original stability.

For a more detailed study, see [8, pág 51] and [14, pág 314]

2.2. Infinite Singular Points. Consider \mathbb{R}^3 the sphere $S^2 = \{(x_1, x_2, x_3) \in \mathbb{R}^3; x_1^2 + x_2^2 + x_3^2 = 1\}$ and the plane $\pi = \{(x_1, x_2, x_3) \in \mathbb{R}^3; x_3 = 1\}$, is tangent to S^2 in the point $(0, 0, 1)$. Let r a line through the origin $(0, 0, 0)$ and a point P of π , then r intercept S^2 in two points P_+ y P_- , where the first is in the upper open hemisphere $H_+ = \{(x_1, x_2, x_3) \in S^2; x_3 > 0\}$ and the second is in the lower open hemisphere $H_- = \{(x_1, x_2, x_3) \in S^2; x_3 < 0\}$.

The expression for $p(X)$ in local chart (U_1, ϕ_1) is given by:

$$(2.5) \quad \begin{cases} \dot{u} &= v^d [-uP(\frac{1}{v}, \frac{u}{v}) + Q(\frac{1}{v}, \frac{u}{v})], \\ \dot{v} &= -v^{d+1}P(\frac{1}{v}, \frac{u}{v}). \end{cases}$$

The expression for (U_2, ϕ_2) is:

$$(2.6) \quad \begin{cases} \dot{u} &= v^d [P(\frac{u}{v}, \frac{1}{v}) - uQ(\frac{u}{v}, \frac{1}{v})], \\ \dot{v} &= -v^{d+1}Q(\frac{u}{v}, \frac{1}{v}). \end{cases}$$

and for (U_3, ϕ_3) is:

$$(2.7) \quad \begin{cases} \dot{u} &= P(u, v), \\ \dot{v} &= Q(u, v). \end{cases}$$

Where d is the maximum degree of the polynomial. For a more detailed study, see [11, pág 151]

2.3. Algebraic Methods. Concerning algebraic aspects considered in this paper, we follow the references [16, 17, 18, 19, 20, 21, 22, 23, 24, 25, 26, 27, 28, 29, 30] and also [1, 2, 3, 4, 5, 6].

A differential field K is a field equiped with a derivation ∂ such that $\forall a, b \in K$ it satisfied:

- (1) $\partial(a + b) = \partial a + \partial b$
- (2) $\partial(a \cdot b) = a \cdot \partial b + \partial a \cdot b$
- (3) $\partial(\frac{a}{b}) = \frac{1}{b^2}(a \cdot \partial b - \partial a \cdot b)$

The field of constants of K , denoted by C_K , is given by

$$C_K = \{c \in K : \partial(c) = 0\}$$

The Picard-Vessiot extension L/K is the extension of K preserving the field of constants, that is $C_L = C_K$. Thus, given a system of first order linear differential equations $\dot{X} = AX$, where $a_{ij} \in K$, the differential Galois group of $\dot{X} = AX$,

denoted by $DGal(L/K)$, is the group of K -differential automorphisms from L to L , i.e., $\sigma : L \mapsto L$, $\partial(\sigma(a)) = \sigma(\partial a)$,

$$DGal(L/K) = \{\sigma : \sigma(a) = a, \forall a \in K\}$$

A Hamiltonian system of n degrees of freedom

$$\dot{q}_i = \frac{\partial H}{\partial p_i}, \quad \dot{p}_i = -\frac{\partial H}{\partial q_i}, \quad 1 \leq i \leq n,$$

where $\mathbf{q} = (q_1, \dots, q_n)$, $\mathbf{p} = (p_1, \dots, p_n)$ and

$$H = \frac{\mathbf{p} \cdot \mathbf{p}}{2m} + V(\mathbf{q}), \quad (\mathbf{q}, \mathbf{p}) \in \mathbb{R}^{2n},$$

is integrable in the Liouville sense whether there exist n independent first integrals that commute in pairwise with the Poisson bracket. In particular, the Hamiltonian systems with one degree of freedom are integrable in the Liouville sense because H is the first integral, i.e., $\dot{H} = 0$. Morales-Ramis theory is the theory that relates differential Galois theory with the integrability of dynamical systems. In particular, Morales-Ramis Theorem for Hamiltonian systems says that *if a Hamiltonian system is integrable, then the connected identity component of the differential Galois group of the first variational equation is an abelian group*.

On the other hand, explicit solutions for differential equation

$$(2.8) \quad \frac{d^2 x}{dt^2} = f(x)$$

are related with the integral curve (x, \dot{x}) of the one degree of freedom Hamiltonian system

$$(2.9) \quad \dot{x} = y, \quad \dot{y} = f(x), \quad H = \frac{y^2}{2} - \int_{x_0}^x f(\tau) d\tau$$

We are interested in the families in where $f(x)$ is a polynomial of degree two, that is, family I, Family IV when $p = -4$ and Family V when $s = -4$. We recall that our problem comes from a polynomial vector field provided in [2, 3], for this reason $p, s \in \mathbb{Z}_0$ to get polynomial vector fields, while $p = -4$ and $s = -4$ correspond originally to a rational non-polynomial vector field, which is exceptionally considered for the algebraic aspects.

Following [31], the Weierstrass P -function is an elliptic function that satisfy the elliptic curve

$$(2.10) \quad y^2 = 4x^3 - g_2x - g_3, \quad x = \wp(t; g_2, g_3), \quad y = \dot{x}.$$

where

$$(2.11) \quad \wp(t; g_2, g_3) = \frac{1}{z^2} + \sum_{\omega} \left(\frac{1}{(t - \omega)^2} - \frac{1}{\omega^2} \right).$$

Moreover, the Weierstrass P -function is a double periodic function with periods $2\omega_1$ and $2\omega_2$; and invariants g_2 and g_3 given by

$$g_2 = \sum_{\omega} \frac{60}{\omega^4}, \quad g_3 = \sum_{\omega} \frac{140}{\omega^6}.$$

The sums range over $\omega = 2n_1\omega_1 + 2n_2\omega_2$ such that $(n_1, n_2) \in \mathbb{Z} \times \mathbb{Z} \setminus \{(0, 0)\}$.

3. Conditions For The Problem.

The following section allows us to identify the quadratic cases associated to (1.1).

3.1. Reduction to 5 Families. The next proposition is a particular case of proposition 4.1 in [3], we only consider the quadratic cases.

PROPOSITION 3.1. *Let $a, b, c, m, k \in \mathbb{R}$ y $s, p, r \in \mathbb{Z}^+$. Quadratic systems associated with each subfamily of (1.1) are equivalently to the following families:*

$$(3.1) \quad \mathbf{I}: \begin{cases} \dot{x} &= y \\ \dot{y} &= -cx^2 \end{cases}$$

$$(3.2) \quad \mathbf{II}: \begin{cases} \dot{x} &= y \\ \dot{y} &= 2byx \end{cases}$$

$$(3.3) \quad \mathbf{III}: \begin{cases} \dot{x} &= y \\ \dot{y} &= 2ayx \end{cases}$$

$$(3.4) \quad \mathbf{IV}: \begin{cases} \dot{x} &= y \\ \dot{y} &= a\left(\frac{p+4}{2}\right)y - \frac{3}{2}a^2x - cx^2 \end{cases}$$

$$(3.5) \quad \mathbf{V}: \begin{cases} \dot{x} &= y \\ \dot{y} &= b\left(\frac{s+4}{2}\right)y - \frac{3}{2}bx - cx^2 \end{cases}$$

PROOF. We analyze each subfamily of the system (1.1), where We observe the different possibilities for the constants a, b and c , are equal to 0. Some of cases are:

I. For $a = 0$, $b = 0$ and $c \neq 0$. We observed two cases:

Case 1. If $s = 0$, then $p = 1$. **Case 2.** If $s = 1$, then $p = 0$.

II. For $a \neq 0$, $b \neq 0$ and $c \neq 0$.

We observed that $\deg(Q) = \max\{2p + 1, 2s + 1, s + p + 1\}$.

Case 1. If $\deg(Q) = 2p + 1$, then $2p + 1 = 2$ so $p = \frac{1}{2} \notin \mathbb{Z}^+$.

Case 2. If $\deg(Q) = 2s + 1$, then $2s + 1 = 2$ so $s = \frac{1}{2} \notin \mathbb{Z}^+$.

Case 3. If $\deg(Q) = s + p + 1$, then we return to reasoning in the family **I**, so $s = 0$ then $p = 1$, but we have that $2p + 1 = 3$ and this case would be cubic. Same for $p = 0$ and $s = 1$. Therefore, this family does not have quadratic cases. □

4. Finite Plane

4.1. Singularity of the Family I.

PROPOSITION 4.1. *(0, 0) is a cusp of family (3.1).*

PROOF. The critical point associated with the system (3.1) is (0, 0). Jacobian matrix is:

$$\mathcal{M}(x, y) = \begin{bmatrix} 0 & 1 \\ -2cx & 0 \end{bmatrix}$$

Then,

$$\mathcal{M}(0,0) = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

We see that $\lambda^2 = 0$. Then, according to the Theorem (2.2), where $A(x, y) = 0$ and $y = 0$, on the other hand We have to $B(x, y) = -cx^2$, we get that $F(x) = -cx^2$ and $G(x) = 0$. Therefore the origin of the system (3.1) is a cusp. \square

4.2. Singularity of the Family II.

PROPOSITION 4.2. The system (3.2) have infinite critical points.

PROOF. $(x, 0)$ which is a line of critical points associated with the system (3.2). We see the solution of the system by separation of variables.

$$y = bx^2 + k, \text{ where } k \text{ is a constant.}$$

\square

4.3. Singularity of the Family III.

PROPOSITION 4.3. The system (3.3) have infinite critical points.

PROOF. $(x, 0)$ which is a line of critical points associated with the system (3.3). We see the solution of the system by separation of variables.

$$y = ax^2 + k, \text{ where } k \text{ is a constant.}$$

\square

4.4. Singularity of the Family IV.

PROPOSITION 4.4. **a)** The point $(0, 0)$ is an stable node if $a < 0$ and unstable if $a > 0$, and $(\frac{-3a^2}{2c}, 0)$ is a saddle.

b) If $p = 0$, $(0, 0)$ is an stable focus if $a < 0$ and unstable if $a > 0$, and $(\frac{-3a^2}{2c}, 0)$ is a saddle.

PROOF. Critical points associated with the system (3.4) are: $(0, 0)$ and $(\frac{-3a^2}{2c}, 0)$. Let $d = a(p + 4)$ and Jacobian matrix are:

$$\mathcal{M}(x, y) = \begin{bmatrix} 0 & 1 \\ -\frac{3}{2}a^2 - 2cx & \frac{d}{2} \end{bmatrix}$$

Then,

$$\mathcal{M}(0,0) = \begin{bmatrix} 0 & 1 \\ -\frac{3}{2}a^2 & \frac{d}{2} \end{bmatrix}$$

and,

$$\mathcal{M}(\frac{-3a^2}{2c}, 0) = \begin{bmatrix} 0 & 1 \\ \frac{3a^2}{2} & \frac{d}{2} \end{bmatrix}$$

a. For $\mathcal{M}(0, 0)$, eigenvalues are:

$\lambda_1 = \frac{1}{4} [d + \sqrt{d^2 - 24a^2}]$ and $\lambda_2 = \frac{1}{4} [d - \sqrt{d^2 - 24a^2}]$. According to the Theorem (2.1) we see that $\lambda_1 \lambda_2 > 0$ therefore $(0, 0)$ is an stable node if $a < 0$ and unstable if $a > 0$.

Now, for $\mathcal{M}(\frac{-3a^2}{2c}, 0)$, eigenvalues are:

$\lambda_1 = \frac{1}{4} [+\sqrt{d^2 + 24a^2}]$ y $\lambda_2 = \frac{1}{4} [d - \sqrt{d^2 + 24a^2}]$. According to the Theorem (2.1) we see that $\lambda_1 \lambda_2 < 0$, then $(\frac{-3a^2}{2c}, 0)$ is a saddle.

b. If $p = 0$, for $\mathcal{M}(0, 0)$ eigenvalues are:

$\lambda_1 = \frac{a}{2} (2 + i\sqrt{2})$ and $\lambda_2 = \frac{a}{2} (2 - i\sqrt{2})$. According to the Theorem (2.1) we see that $(0, 0)$ is an stable focus if $a < 0$ and unstable if $a > 0$.

Now, for $\mathcal{M}(\frac{-3a^2}{2c}, 0)$, eigenvalues are:

$\lambda_1 = \frac{a}{4} [4 + 2\sqrt{10}]$ and $\lambda_2 = \frac{a}{4} [4 - 2\sqrt{10}]$. According to the Theorem (2.1) we see that $\lambda_1 \lambda_2 < 0$, then $(\frac{-3a^2}{2c}, 0)$ is a saddle.

□

4.5. Singularity of the Family V. Before looking at the following proposition, We define the following regions:

$$\begin{aligned} R_1 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b > 0\} \\ R_2 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b = 0\} \\ R_3 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b < 0, c > 0\} \\ R_4 &= \{(b, 0, d) \in \mathbb{R}^3 | d^2 - 24b > 0\} \\ R_5 &= \{(b, 0, d) \in \mathbb{R}^3 | d^2 - 24b = 0\} \\ R_6 &= \{(b, 0, d) \in \mathbb{R}^3 | d^2 - 24b < 0\} \\ R_7 &= \{(0, c, 0) \in \mathbb{R}^3 | c > 0\} \\ R_8 &= \{(b, c, d) \in \mathbb{R}^3 | c < 0\} \end{aligned}$$

We note that $\mathbb{R}^3 = \bigcup_{i=1}^8 R_i$. Now in R_3 and R_4 we consider the following subsets:

$$\begin{aligned} E_1 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b < 0, d > 0, c > 0\} \\ E_2 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b < 0, d < 0, c > 0\} \\ E_3 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 + 24b < 0, d > 0, c > 0\} \\ E_4 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b < 0, d < 0, c > 0\} \\ E_5 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b < 0, d > 0, c < 0\} \\ E_6 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b < 0, d < 0, c < 0\} \\ E_7 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 + 24b < 0, d > 0, c < 0\} \\ E_8 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b < 0, d < 0, c < 0\} \\ E_9 &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b > 0, d > 0, c > 0\} \\ E_{10} &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b > 0, d < 0, c > 0\} \\ E_{11} &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b > 0, d < 0, c < 0\} \\ E_{12} &= \{(b, c, d) \in \mathbb{R}^3 | d^2 - 24b > 0, d > 0, c < 0\} \end{aligned}$$

PROPOSITION 4.5. Let the family (3.5) with $(b, c, d) \in \mathbb{R}^3$, then:

- a) If $(b, c, d) \in R_1$ and $b > 0$ then the point $(0, 0)$ is unstable node and the point $(\frac{-3b}{2c}, 0)$ is a saddle. if $b < 0$ then the point $(0, 0)$ is a saddle and the point $(\frac{-3b}{2c}, 0)$ is stable node.
- b) If $(b, c, d) \in R_2$ and $b > 0$, then the critical point $(0, 0)$ is a unstable node and the critical point $(\frac{-3b}{2c}, 0)$ is saddle.
- c) If $(b, c, d) \in R_3$ and $b > 0$ then the point $(0, 0)$ is stable focus and the point $(\frac{-3b}{2c}, 0)$ is a saddle. If $b < 0$ then point $(0, 0)$ is a unstable focus and the point $(\frac{-3b}{2c}, 0)$ is a unstable node.

PROOF. Let $d = b(s + 4)$, so critical points associated with the system (3.5) are: $(0, 0)$ and $(\frac{-3b}{2c}, 0)$.
then, Jacobian matrix are:

$$\mathcal{M}(x, y) = \begin{bmatrix} 0 & 1 \\ -\frac{3}{2}b - 2cx & \frac{d}{2} \end{bmatrix}$$

Then,

$$\mathcal{M}(0, 0) = \begin{bmatrix} 0 & 1 \\ -\frac{3b}{2} & \frac{d}{2} \end{bmatrix}$$

For $\mathcal{M}(0, 0)$, eigenvalues are:

$$\lambda_1 = \frac{1}{4} [d + \sqrt{d^2 - 24b}] \text{ and } \lambda_2 = \frac{1}{4} [d - \sqrt{d^2 - 24b}].$$

Now, for $(\frac{-3b}{2c}, 0)$, we have that

$$\mathcal{M}(\frac{-3b}{2c}, 0) = \begin{bmatrix} 0 & 1 \\ \frac{3b}{2} & \frac{d}{2} \end{bmatrix}$$

With eigenvalues:

$$\lambda_1 = \frac{1}{4} [d + \sqrt{d^2 + 24b}] \text{ and } \lambda_2 = \frac{1}{4} [d - \sqrt{d^2 + 24b}].$$

- a. If $(b, c, d) \in R_1$ that is $d^2 - 24b > 0$, for $\mathcal{M}(0, 0)$ We see that $\lambda_1 \lambda_2 = \frac{-3b}{2}$ and $\lambda_1 > 0$, then according to the Theorem (2.1) the critical point $(0, 0)$ is unstable node if $b > 0$ and a saddle if $b < 0$.

Now, for $\mathcal{M}(\frac{-3b}{2c}, 0)$, we have that $d^2 + 24b > 48b$, then:

- (1) If $b > 0$, According to the Theorem (2.1) We see that $\lambda_1 \lambda_2 = \frac{-3b}{2}$, and $b > 0$ then $(\frac{-3b}{2c}, 0)$ is a saddle. If $b < 0$, $\lambda_2 < 0$, then stable node.
- (2) If $b < 0$ and $d^2 + 24b \in [48b, 0)$, then the critical point $(\frac{-3b}{2c}, 0)$ is stable focus.

(3) If $b < 0$ and $d^2 + 24b \geq 0$, then the critical point $(\frac{-3b}{2c}, 0)$ is a stable node.

b. If $(b, c, d) \in R_2$ that is $d^2 - 24b = 0$ this leads to $b \geq 0$ for $\mathcal{M}(0, 0)$ $\lambda_1 = \lambda_2 = \frac{d}{4}$, if $b > 0$ then the critical point is an unstable node. We note that if $b = 0$, then this corresponds to (3.1). Now, for $\mathcal{M}(\frac{-3b}{2c}, 0)$, if $b > 0$ that is $d^2 - 24b > 0$, Furthermore $\lambda_1 \lambda_2 = \frac{-3b}{2}$ according to the Theorem (2.1) the critical point $(\frac{-3b}{2c}, 0)$ is a saddle.

c. If $(b, c, d) \in R_3$, that is $d^2 - 24b < 0$ then $b > 0$. For $\mathcal{M}(0, 0)$ eigenvalues are: $\lambda_1 = \frac{1}{4}(d + i\sqrt{24b - d^2})$ and $\lambda_2 = \frac{1}{4}(d - i\sqrt{24b - d^2})$, then according to the Theorem (2.1), $(0, 0)$ is a focus unstable. Now, for $b > 0$ We see that $d^2 + 24b > 0$, that is for $\mathcal{M}(\frac{-3b}{2c}, 0)$, we have that $\lambda_{1,2} \in \mathbb{R}$. So, $\lambda_1 \lambda_2 = \frac{-3b}{2}$, so According to the Theorem (2.1) we have that the critical point $(\frac{-3b}{2c}, 0)$ is a saddle. □

PROPOSITION 4.6. Given the family (3.5) with $c = 0$, then:

- a)** If $(b, 0, d) \in R_4$ and $b > 0$, then the critical point $(0, 0)$ is a saddle. If $b < 0$, then the critical point $(0, 0)$ is a stable node.
- b)** If $(b, 0, d) \in R_5$ and $b > 0$, then the critical point $(0, 0)$ is an unstable node.
- c)** If $(b, 0, d) \in R_6$ and $b > 0$, then the critical point $(0, 0)$ is an unstable focus. If $b < 0$, then the critical point $(0, 0)$ is a stable focus.

PROOF. With $c = 0$, the family (3.5) have the form:

$$(4.1) \quad \begin{cases} \dot{x} &= y \\ \dot{y} &= \frac{d}{2}y - \frac{3}{2}bx \end{cases}$$

Here, we see that the only critical point associated with the family (4.1) is $(0, 0)$ then, Jacobian matrix evaluated in the point is:

$$\mathcal{M}(0, 0) = \begin{bmatrix} 0 & 1 \\ -\frac{3b}{2} & \frac{d}{2} \end{bmatrix}$$

For $\mathcal{M}(0, 0)$, eigenvalues are:

$$\lambda_1 = \frac{1}{4} [d + \sqrt{d^2 - 24b}] \text{ and } \lambda_2 = \frac{1}{4} [d - \sqrt{d^2 - 24b}].$$

- a.** If $(b, 0, d) \in R_4$ that is $d^2 - 24b > 0$, we have that $\lambda_1 \lambda_2 = \frac{-3b}{2}$, then if $b > 0$ according to the Theorem (2.1) the critical point $(0, 0)$ is a saddle and if $b < 0$ and $\lambda_1 < 0$ then the critical point $(0, 0)$ is a stable node.
- b.** If $(b, 0, d) \in R_5$ that is $d^2 - 24b = 0$, we have that $\lambda_1 \lambda_2 = \frac{d}{4}$, then according to the Theorem (2.1) the critical point $(0, 0)$ is an unstable node.

- c. If $(b, 0, d) \in R_6$ that is $d^2 - 24b < 0$, we have that $\lambda_1 = \frac{1}{4}(d + i\sqrt{24b - d^2})$ and $\lambda_2 = \frac{1}{4}(d - i\sqrt{24b - d^2})$, then if $b > 0$ according to the Theorem (2.1) the critical point $(0, 0)$ is a unstable focus and if $b < 0$ and $\lambda_1 < 0$ then the critical point $(0, 0)$ is a stable focus.

□

Now, We will look for the stable manifold associated systems.

PROPOSITION 4.7. The stable manifold associated with the system (3.4) at the point $(\frac{-3a^2}{2c}, 0)$ is:

$$S : y = \frac{c(x + \frac{3a^2}{2c})^2}{(v-w)(v-2w)}$$

PROOF. Let is observe the stability of the system (3.4) in the point $(\frac{-3a^2}{2c}, 0)$: Let is look at the eigenvalues in the Jacobian matrix of the system (3.4) in the point $(\frac{-3a^2}{2c}, 0)$.

$$A = \begin{bmatrix} 0 & 1 \\ \frac{3a^2}{2} & d \end{bmatrix}$$

So, $w = \lambda_1 = \frac{1}{4}[d + \sqrt{d^2 + 24a^2}]$ and $v = \lambda_2 = \frac{1}{4}[d - \sqrt{d^2 + 24a^2}]$.

Then, $B(x) = C^{-1}AC = \begin{bmatrix} w & 0 \\ 0 & v \end{bmatrix}$

$F(x) = \begin{bmatrix} 0 \\ -cx^2 \end{bmatrix}$, $G(x) = \frac{cx^2}{v-w} \begin{bmatrix} 1 \\ -1 \end{bmatrix}$

$$U(t) = \begin{bmatrix} e^{wt} & 0 \\ 0 & 0 \end{bmatrix}, \quad V(t) = \begin{bmatrix} 0 & 0 \\ 0 & e^{vt} \end{bmatrix}, \quad b = \begin{bmatrix} b_1 \\ 0 \end{bmatrix}$$

Then,

$u^{(0)}(t, b) = 0$.

$u^{(1)}(t, b) = \begin{bmatrix} e^{wt}b_1 \\ 0 \end{bmatrix}$

$u^{(2)}(t, b) = \begin{bmatrix} e^{wt}b_1 \\ 0 \end{bmatrix} + \int_0^t \begin{bmatrix} e^{w(t-s)} & 0 \\ 0 & 0 \end{bmatrix} \left[\begin{bmatrix} \frac{b_1^2c}{v-w}e^{2ws} \\ -\frac{b_1^2c}{v-w}e^{2ws} \end{bmatrix} ds - \int_t^\infty \begin{bmatrix} 0 & 0 \\ 0 & e^{v(t-s)} \end{bmatrix} \left[\begin{bmatrix} \frac{b_1^2c}{v-w}e^{2ws} \\ -\frac{b_1^2c}{v-w}e^{2ws} \end{bmatrix} ds =$

$u^{(2)}(t, b) = \begin{bmatrix} e^{wt}b_1 + \frac{b_1^2ce^{2wt}[e^{wt}-1]}{w(v-w)} \\ \frac{b_1^2ce^{2wt}}{(v-w)(v-2w)} \end{bmatrix}$

Therefore, We can approximate by $\psi_2(b_1) = b_1$, therefore the stable manifold can be approximated by

$$S : y = \frac{c(x + \frac{3a^2}{2c})^2}{(v-w)(v-2w)}$$

like $x \rightarrow 0$. Similarly the unstable

$$U : x + \frac{3a^2}{2c} = \frac{2cy^2}{(v-w)(v-2w)}$$

□

PROPOSITION 4.8. For the system (3.5) We have that:

a) If $(b, c, d) \in R_1$ and $b < 0$, stable manifold at the point $(0, 0)$ is:

$$S : y = \frac{cx^2}{(v-w)(v-2w)}$$

b) If $(b, c, d) \in \{(x, y, z)/x > 0, y \neq 0\}$, then stable manifold at the point $(\frac{-3b}{2c}, 0)$ es:

$$S : y = \frac{c(x + \frac{3b}{2c})^2}{(v-w)(v-2w)}$$

PROOF. a) Let is observe the stability of the system (3.5) for $b < 0$ at the point $(0, 0)$:

Let is observe the stability of the system (3.5) in the point $(0, 0)$.

$$A = \begin{bmatrix} 0 & 1 \\ \frac{-3b}{2} & \frac{d}{2} \end{bmatrix}$$

$$\text{Let, } w = \lambda_1 = \frac{1}{4}[d + \sqrt{d^2 - 24b}] \quad \text{and} \quad v = \lambda_2 = \frac{1}{4}[d - \sqrt{d^2 - 24b}].$$

$$\text{Then, } B(x) = C^{-1}AC = \begin{bmatrix} w & 0 \\ 0 & v \end{bmatrix}$$

$$F(x) = \begin{bmatrix} 0 \\ -cx^2 \end{bmatrix}, \quad G(x) = \frac{cx^2}{v-w} \begin{bmatrix} 1 \\ -1 \end{bmatrix}$$

$$U(t) = \begin{bmatrix} e^{wt} & 0 \\ 0 & 0 \end{bmatrix}, \quad V(t) = \begin{bmatrix} 0 & 0 \\ 0 & e^{vt} \end{bmatrix}, \quad a = \begin{bmatrix} a_1 \\ 0 \end{bmatrix}$$

Then,

$$u^{(0)}(t, a) = 0.$$

$$u^{(1)}(t, a) = \begin{bmatrix} e^{wt}a_1 \\ 0 \end{bmatrix}$$

$$u^{(2)}(t, a) = \begin{bmatrix} e^{wt}a_1 \\ 0 \end{bmatrix} + \int_0^t \begin{bmatrix} e^{w(t-s)} & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} \frac{a_1^2c}{v-w}e^{2ws} \\ -\frac{a_1^2c}{v-w}e^{2ws} \end{bmatrix} ds - \int_t^\infty \begin{bmatrix} 0 & 0 \\ 0 & e^{v(t-s)} \end{bmatrix} \begin{bmatrix} \frac{a_1^2c}{v-w}e^{2ws} \\ -\frac{a_1^2c}{v-w}e^{2ws} \end{bmatrix} ds =$$

$$u^{(2)}(t, a) = \begin{bmatrix} e^{wt}a_1 + \frac{a_1^2ce^{2wt}[e^{wt}-1]}{w(v-w)} \\ \frac{a_1^2ce^{2wt}}{(v-w)(v-2w)} \end{bmatrix}$$

Therefore, We can approximate by $\psi_2(a_1) = a_1$, therefore the stable manifold can be approximated by

$$S : y = \frac{cx^2}{(v-w)(v-2w)}$$

Like $x \rightarrow 0$. Similarly the unstable

$$U : x = \frac{2cy^2}{(v-w)(v-2w)}$$

b) Let is observe the stability of the system (3.5) at the point $(\frac{-3b}{2c}, 0)$, when $b > 0$:

$$A = \begin{bmatrix} 0 & 1 \\ \frac{3b}{2} & \frac{d}{2} \end{bmatrix}$$

Let, $w = \lambda_1 = \frac{1}{4}[d + \sqrt{d^2 + 24b}]$ and $v = \lambda_2 = \frac{1}{4}[d - \sqrt{d^2 + 24b}]$.

$$\text{So, } B(x) = C^{-1}AC = \begin{bmatrix} w & 0 \\ 0 & v \end{bmatrix}$$

$$F(x) = \begin{bmatrix} 0 \\ -cx^2 \end{bmatrix}, \quad G(x) = \frac{cx^2}{v-w} \begin{bmatrix} 1 \\ -1 \end{bmatrix}$$

$$U(t) = \begin{bmatrix} e^{wt} & 0 \\ 0 & 0 \end{bmatrix}, \quad V(t) = \begin{bmatrix} 0 & 0 \\ 0 & e^{vt} \end{bmatrix}, \quad a = \begin{bmatrix} a_1 \\ 0 \end{bmatrix}$$

Then,

$$u^{(0)}(t, a) = 0.$$

$$u^{(1)}(t, a) = \begin{bmatrix} e^{wt} a_1 \\ 0 \end{bmatrix}$$

$$u^{(2)}(t, a) = \begin{bmatrix} e^{wt} a_1 \\ 0 \end{bmatrix} + \int_0^t \begin{bmatrix} e^{w(t-s)} & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} \frac{a_1^2 c}{v-w} e^{2ws} \\ -\frac{a_1^2 c}{v-w} e^{2ws} \end{bmatrix} ds - \int_t^\infty \begin{bmatrix} 0 & 0 \\ 0 & e^{v(t-s)} \end{bmatrix} \begin{bmatrix} \frac{a_1^2 c}{v-w} e^{2ws} \\ -\frac{a_1^2 c}{v-w} e^{2ws} \end{bmatrix} ds =$$

$$u^{(2)}(t, a) = \begin{bmatrix} e^{wt} a_1 + \frac{a_1^2 c e^{2wt} [e^{wt} - 1]}{w(v-w)} \\ \frac{a_1^2 c e^{2wt}}{(v-w)(v-2w)} \end{bmatrix}$$

Therefore, We can approximate by $\psi_2(a_1) = a_1$, therefore the stable manifold can be approximated by

$$S : y = \frac{c(x + \frac{3b}{2c})^2}{(v-w)(v-2w)}$$

Like $x \rightarrow 0$. Similarly the unstable

$$U : x + \frac{3b}{2c} = \frac{2cy^2}{(v-w)(v-2w)}$$

□

5. Bifurcations

In this section we will analyze the study of the bifurcations of family (3.5)

5.1. Family V.

PROPOSITION 5.1. *Let sets R_7 and R_8 are transcritical bifurcations for the system (3.5)*

PROOF. Let $P_1 : (0, 0)$ and $P_2 : (\frac{3b}{2c}, 0)$ of proposition (4.5). If $(b, c, d) \in E_3$ then P_1 a saddle and P_2 is a unstable focus, when $(b, c, d) \in R_7$, P_1 and P_2 , they collapse on one critical point which point is a cusp. So, when $(b, c, d) \in E_2$ then P_1 a unstable focus and P_2 is a saddle. Similarly, the same behavior is observed when $(b, c, d) \in E_2$, then $(b, c, d) \in R_7$ and finally $(b, c, d) \in E_4$.

Now, Let $P_1 : (0, 0)$ and $P_2 : (\frac{3b}{2c}, 0)$ of proposition (4.5). If $(b, c, d) \in E_4$ then P_1 a saddle and P_2 is a stable focus, when $(b, c, d) \in R_8$, P_1 and P_2 , they collapse on one critical point which point is a cusp. So, when $(b, c, d) \in E_1$ then P_1 a stable focus and P_2 is a saddle. Similarly, the same behavior is observed when $(b, c, d) \in E_1$, then $(b, c, d) \in R_8$ and finally $(b, c, d) \in E_3$.

Therefore, sets R_7 and R_8 are transcritical bifurcations for the system (3.5) \square

PROPOSITION 5.2. A set $\{(b, 0, d) | d^2 - 24b < 0\}$ is a bifurcations saddle-focus-saddle for the system (3.5).

PROOF. For proposition (4.5), if $(b, c, d) \in E_1$, the point P_1 is a stable focus and P_2 is a saddle. Now, when $(b, c, d) \in \{(b, 0, d) | d^2 - 24b < 0, d > 0\}$ for proposition (4.5), P_1 and P_2 they collapse in an unstable focus when (b, c, d) goes the set E_5 , appear again P_1 and P_2 like stable focus and a saddle respectively. \square

PROPOSITION 5.3. A set $\{(b, 0, d) | d^2 - 24b < 0\}$ is a bifurcations saddle-focus-saddle for the system (3.5).

PROOF. For proposition (4.5), if $(b, c, d) \in E_2$, the point P_1 is a unstable focus and P_2 is a saddle. Now, when $(b, c, d) \in \{(b, 0, d) | d^2 - 24b < 0, d < 0\}$ for proposition (4.5), P_1 and P_2 they collapse in an stable focus when (b, c, d) goes the set E_6 , appear again P_1 and P_2 like unstable focus and a saddle respectively. \square

PROPOSITION 5.4. Let sets E_9 and E_{10} are local bifurcations for the system (3.5)

PROOF. Let $P_1 : (0, 0)$ of proposition (4.5). If $(b, c, d) \in E_9$ then P_1 is a stable node. Now, when $(b, c, d) \in E_{10}$ the point P_1 is a unstable node. Therefore, regions E_9 and E_{10} are local bifurcations for the system (3.5) \square

PROPOSITION 5.5. Let sets E_{11} and E_{12} are local bifurcations for the system (3.5)

PROOF. Let $P_2 : (\frac{3b}{2c}, 0)$ of proposition (4.5). If $(b, c, d) \in E_{12}$ then P_2 is a unstable node. Now, when $(b, c, d) \in E_{11}$ the point P_2 is a stable node. Therefore, regions E_{11} and E_{12} are local bifurcations for the system (3.5) \square

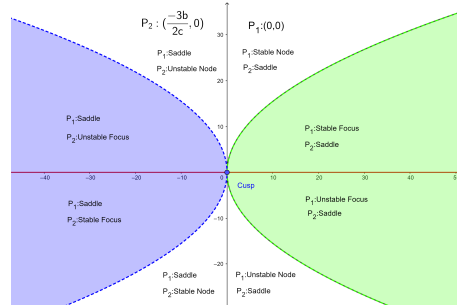


FIGURE 1. (3.5), $c > 0$.

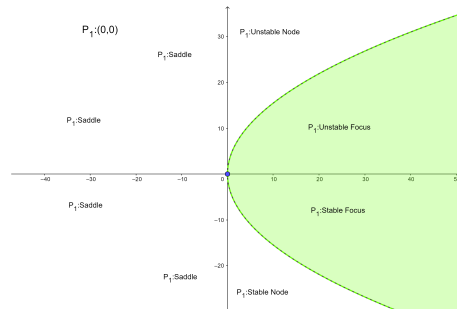


FIGURE 2. (3.5), $c = 0$.

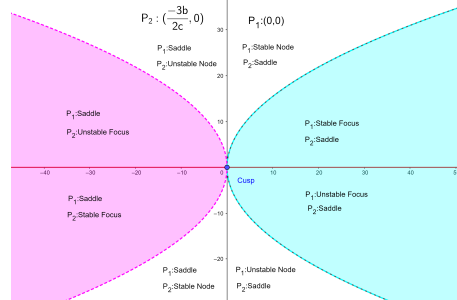


FIGURE 3. (3.5), $c < 0$.

6. Infinite Plane

6.1. Family I. In the Chart U_1 the associated system 3.1 :

$$(6.1) \quad \begin{cases} \dot{u} = -u^2v - c \\ \dot{v} = -uv^2 \end{cases}$$

The system have not critical points at infinite plane.

In the Chart U_2 the associated system 3.1:

$$(6.2) \quad \begin{cases} \dot{u} = v + cu^3 \\ \dot{v} = -cu^2v \end{cases}$$

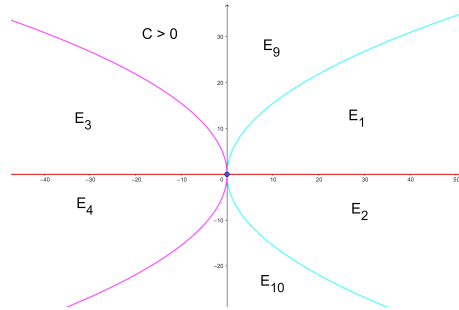


FIGURE 4. (3.5), $c > 0$.

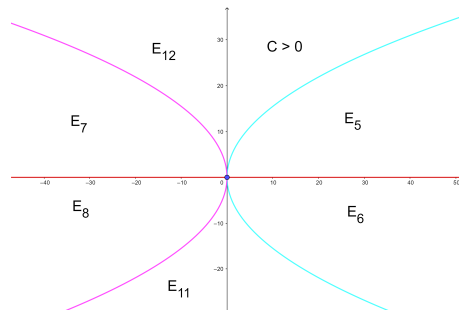


FIGURE 5. (3.5), $c < 0$.

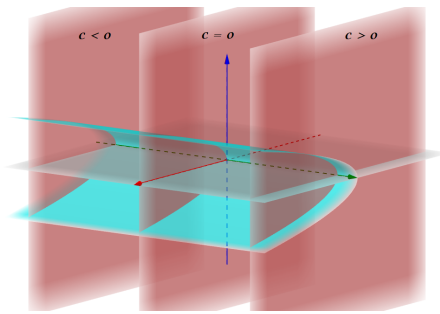


FIGURE 6. (3.5), $c < 0$.

PROPOSITION 6.1. *The point $(0, 0)$ is an stable node if $c < 0$ and unstable if $c > 0$.*

PROOF. The critical points associated with the system (6.2) is $P : (0, 0)$. Jacobian matrix is:

$$\mathcal{M}(u, v) = \begin{bmatrix} 3cu^2 & 1 \\ -2cuv & -cu^2 \end{bmatrix}$$

Then,

$$\mathcal{M}(0, 0) = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

We see that $\lambda^2 = 0$. According to the Theorem (2.2), let $v + A(u, v) = 0$ a solution of $v + A(u, v) = 0$, where $A(u, v) = cu^3$ then $v = -cu^3$, also we have that $B(u, v) = cu^2v$, so $F(u) = -c^2u^5$ and $G(x) = 4cu^2$ then $m = 5$, $n = 2$, $a = -c^2$, $b = 4c$ and $m = 2n + 1$, furthermore $b^2 + 4a(n + 1) \geq 0$. Therefore the origin of the system (6.2) in infinite plane is an stable node if $c < 0$ and unstable if $c > 0$. \square

For a more detailed study of the system see figure on the Poincaré sphere (1).

6.2. Family II. In the Chart U_1 the associated system (3.2):

$$(6.3) \quad \begin{cases} \dot{u} &= -u^2v + 2b \\ \dot{v} &= -uv^2 \end{cases}$$

The system have not critical points at infinite plane.

In the Chart U_2 the associated system (3.2).

$$(6.4) \quad \begin{cases} \dot{u} &= v - 2bu^2 \\ \dot{v} &= -2buv \end{cases}$$

PROPOSITION 6.2. The point $(0, 0)$ have one hyperbolic and one elliptic sector.

PROOF. The critical points associated with the system (6.4) is $P : (0, 0)$.
Jacobian matrix is:

$$\mathcal{M}(u, v) = \begin{bmatrix} -4bu & 1 \\ -2bv & -2bu \end{bmatrix}$$

Then,

$$\mathcal{M}(0, 0) = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

We see that, $\lambda^2 = 0$. According to the Theorem (2.2), let $v + A(u, v) = 0$ a solution of $v + A(u, v) = 0$, where $A(u, v) = -2bu^2$ then $v = -2bu^2$, also we have that $B(u, v) = -2buv$, so $F(u) = -4b^2u^3$ and $G(x) = -6bu$ then $m = 2n + 1$ and $b^2 + 4a(n + 1)$. Therefore the origin of the system (6.4) in infinite plane have one hyperbolic and one elliptic sector. \square

For a more detailed study of the system see figure on the Poincaré sphere (2).

6.3. Family III. In the Chart U_1 the associated system (3.3):

$$(6.5) \quad \begin{cases} \dot{u} &= -u^2v + 2a \\ \dot{v} &= -uv^2 \end{cases}$$

The system have not critical points at infinite plane.

In the Chart U_2 the associated system (3.3).

$$(6.6) \quad \begin{cases} \dot{u} &= v - 2au^2 \\ \dot{v} &= -2auv \end{cases}$$

PROPOSITION 6.3. The point $(0, 0)$ have one hyperbolic and one elliptic sector.

PROOF. The critical points associated with the system (6.6) is $P : (0, 0)$.
Jacobian matrix is:

$$\mathcal{M}(u, v) = \begin{bmatrix} -4au & 1 \\ -2av & -2au \end{bmatrix}$$

Then,

$$\mathcal{M}(0,0) = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

We see that, $\lambda^2 = 0$. According to the Theorem (2.2), let $v + A(u, v) = 0$ a solution of $v + A(u, v) = 0$, where $A(u, v) = -2au^2$ then $v = -2au^2$, also we have that $B(u, v) = -2auv$, so $F(u) = -4a^2u^3$ y $G(x) = -6au$ then $m = 2n + 1$ and $b^2 + 4a(n + 1)$. Therefore the origin of the system (6.6) in infinite plane have one hyperbolic and one elliptic sector. \square

For a more detailed study of the system see figure on the Poincaré sphere (3).

6.4. Family IV. Let $d = a(p + 4)$.

In the Chart U_1 the associated system (3.4).

$$(6.7) \quad \begin{cases} \dot{u} &= -u^2v + \frac{dvw}{2} - \frac{3a^2v}{2} - c \\ \dot{v} &= -uv^2 \end{cases}$$

The system have not critical points at infinite plane.

In the Chart U_2 the associated system (3.4).

$$(6.8) \quad \begin{cases} \dot{u} &= v - \frac{dvw}{2} + \frac{3}{2}a^2u^2v + cu^3 \\ \dot{v} &= -\frac{dv^2}{2} + \frac{3}{2}a^2uv^2 + cu^2v \end{cases}$$

PROPOSITION 6.4. The point $(0,0)$ is an stable node if $c < 0$ and unstable if $c > 0$.

PROOF. The critical points associated with the system (6.8) is $P : (0,0)$.

Jacobian matrix:

$$\mathcal{M}(u,v) = \begin{bmatrix} -\frac{dv}{2} + 3a^2v + 3cu^2 & 1 - \frac{du}{2} + \frac{3}{2}a^2u^2 \\ \frac{3}{2}a^2v^2 & -dv + 3a^2uv + cu^2 \end{bmatrix}$$

Then,

$$\mathcal{M}(0,0) = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

We see that, $\lambda^2 = 0$. According to the Theorem (2.2), Let $v = f(u)$ a solution of $v + A(u, v) = 0$ where $v = f(u) = -cu^3 + \dots$ an approximation of the Taylor series solution, furthermore $B(u, v) = -\frac{dv^2}{2} + \frac{3}{2}a^2uv^2 + cu^2v$, then $F(u) = -c^2u^5 + \dots$ and $G(u) = cu^2 + \dots$, so $m = 5, n = 2, b = 4c$ and $a = -c^2$. Therefore the origin in the infinite plane is an stable node if $c < 0$ and unstable if $c > 0$. \square

For a more detailed study of the system see figure on the Poincaré sphere (4).

6.5. Family V. Let $d = a(s + 4)$.

In the Chart U_1 the associated system (3.5).

$$(6.9) \quad \begin{cases} \dot{u} &= -u^2v + \frac{duv}{2} - \frac{3bv}{2} - c \\ \dot{v} &= -uv^2 \end{cases}$$

The system have not critical points at infinite plane.

In the Chart U_2 the associated system (3.5)

$$(6.10) \quad \begin{cases} \dot{u} &= v - \frac{duv}{2} + \frac{3}{2}bu^2v + cu^3 \\ \dot{v} &= -\frac{dv^2}{2} + \frac{3}{2}buv^2 + cu^2v \end{cases}$$

PROPOSITION 6.5. The point $(0, 0)$ is stable node if $c < 0$ and unstable if $c > 0$.

PROOF. The critical points associated with the system (6.10) is $P : (0, 0)$.

Jacobian matrix:

$$\mathcal{M}(u, v) = \begin{bmatrix} -\frac{dv}{2} + 3buv + 3cu^2 & 1 - \frac{du}{2} + \frac{3}{2}bu^2 \\ \frac{3}{2}bv^2 + 2cuv & -2dv + 3buv + cu^2 \end{bmatrix}$$

Then,

$$\mathcal{M}(0, 0) = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$$

We see that, $\lambda^2 = 0$. According to the Theorem (2.2), let $v = f(u)$ a solution of $v + A(u, v) = 0$, where $v = f(u) = -cu^3 + \dots$ approximation of the Taylor series solution, furthermore $B(u, v) = -\frac{dv^2}{2} + \frac{3}{2}buv^2 + cu^2v$, then $F(u) = -c^2u^5 + \dots$ y $G(u) = cu^2 + \dots$, so $m = 5, n = 2, b = 4c$ y $a = -c^2$. Therefore the origin in the infinite plane is an stable node if $c < 0$ and is a unstable node if $c > 0$. \square

For a more detailed study of the system which can see figure on the Poincarè sphere (5) and (6) .

7. Global Phase Portrait

In this section We show the global phase portrait associate to each family:

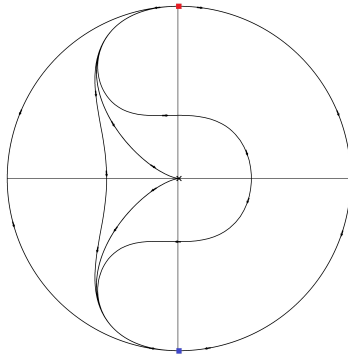


FIGURE 7. Family I

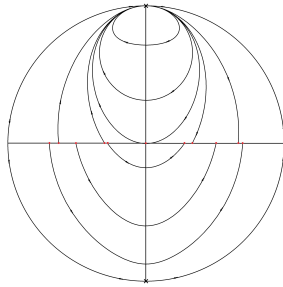


FIGURE 8. Family II

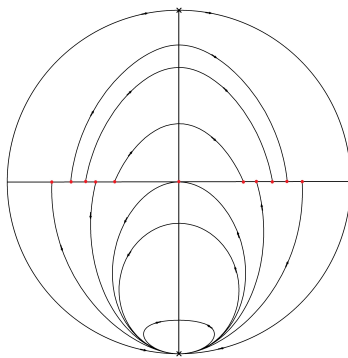


FIGURE 9. Family III

8. Algebraic Aspects

In this section we analyze the families I, II, III, IV and V through an algebraic point of view. We compute the solutions in terms of P-Weierstrass function of families I, II ($p = -4$) and V ($s = -4$), as well the differential Galois group of their variational equations.

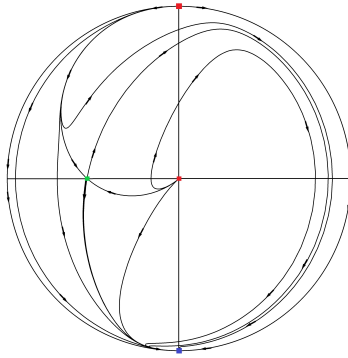


FIGURE 10. Family IV

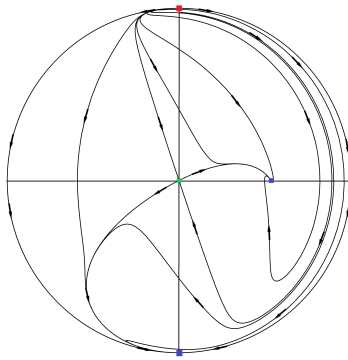


FIGURE 11. Family V when $b < 0$.

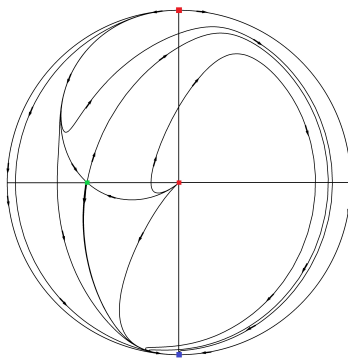


FIGURE 12. Family V when $b > 0$.

8.1. Family I.

THEOREM 8.1. *Consider the family I, The following statements hold.*

- (1) *The dynamical system is hamiltonian with one degree of freedom and with polynomial first integral*

$$H = H(x, y) = \frac{y^2}{2} + \frac{c}{3}x^3$$

- (2) *The integral curve of the Hamiltonian vector field is*

$$\left(-\frac{6}{c}\wp(t + k_0; 0, -2H), -\frac{6}{c}\dot{\wp}(t + k_0; 0, -2H)\right).$$

- (3) *The Differential Galois Group associated to the foliation is isomorphic to \mathbb{Z}_2 .*
 (4) *The connected identity component of the Differential Galois Group of the first variational equation along any particular solution is an abelian group.*

PROOF. We proceed according to each item,

- (1) The polynomial vector field related with family I is equivalent to Equation (2.8) being $f(x) = -cx^2$. In virtue of Equation (2.9) we have the Hamiltonian $H = \frac{y^2}{2} + \frac{c}{3}x^3$.
 (2) Due to $y = \dot{x}$, we obtain $y^2 = -\frac{2c^2}{3} + H$. Through the change of variable $(x, y) \mapsto (\sqrt[3]{\frac{-6}{c}}x, \sqrt[3]{\frac{-6}{c}}y)$, we arrive to the elliptic curve given in Equation (2.10) with invariants $g_2 = 0$ and $g_3 = -2H$. Thus, the integral curve of the Hamiltonian system is (x, \dot{x}) , being x given by $-\frac{6}{c}\wp(t + k_0; 0, -2H)$.
 (3) The foliation associated to the vector field of Family I is

$$y' = -\frac{cx^2}{y}, \quad ' := \frac{d}{dx}.$$

Setting $z = \frac{y^2}{2}$, we obtain $z' = -cx^2$ and therefore $z = -\frac{c}{3}x^3$. Due to the differential field K is the field of rational functions $\mathbb{C}(x)$, $\sigma(z) = z$ and $\sigma(y) = \lambda\sqrt{z}$, where $\lambda^2 = 1$. Thus, the Picard-Vessiot extension L is a quadratic extension of K and we can conclude that $DGal(L/K)$ has two elements.

- (4) Let $(x_0(t), \dot{x}_0(t))$ be a particular solution of the polynomial vector field related with Family I. Thus, the first variational equation is

$$\frac{d}{dt} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -2cx_0(t) & 0 \end{pmatrix} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix},$$

which is equivalent to $\ddot{\xi} = -2cx_0(t)\xi$, being $\xi = \xi_1$. By Morales-Ramis theory, due to the dynamical system is polynomially integrable, the differential Galois group of this first variational equation is abelian. □

8.2. Family II.

THEOREM 8.2. *Consider the family II, The following statements hold.*

- (1) *The first integral of the polynomial vector field is*

$$I = I(x, y) = y - bx^2$$

- (2) *The integral curve of the polynomial vector field is $(x(t), \dot{x}(t))$, where*

$$x(t) = \sqrt{\frac{k_1}{b}} \tan(\sqrt{k_1 b(k_2 + t)}).$$

- (3) *The Differential Galois Group associated to the foliation is isomorphic to the identity group.*
- (4) *The connected identity component of the Differential Galois Group of the first variational equation around any particular solution is an abelian group.*

PROOF. We proceed according to each item,

- (1) The total derivative of $I(x, y)$ vanishes, i.e., $\dot{I} = 0$, therefore I is a first integral of the vector field related to family II.
- (2) Due to $y = \dot{x}$, we obtain $\ddot{x} = b\dot{z}$, where $z = x^2$. Thus, $\dot{x} = bx^2 + k_1$, which implies that

$$\int \frac{dx}{bx^2 + k_1} = t + k_2$$

and then $x(t) = \sqrt{\frac{k_1}{b}} \tan(\sqrt{k_1 b}(k_2 + t))$.

- (3) The foliation associated to the vector field of Family II is

$$y' = 2bx, \quad ' := \frac{d}{dx}.$$

Then the solution of this foliation is

$$y(x) = bx^2 + k_1$$

. Then we can conclude that $DGal(L/K)$ has one element, i.e., $DGal(L/K) = I_2$.

- (4) Let $(x_0(t), \dot{x}_0(t))$ be a particular solution of the polynomial vector field related with Family II. Thus, the first variational equation is

$$\frac{d}{dt} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -2by_0(t) & 2bx_0(t) \end{pmatrix} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix},$$

which is equivalent to

$$\ddot{\xi} - 2bx_0(t)\dot{\xi} - 2by_0(t)\xi = 0, \quad \xi = \xi_1.$$

Due to the first integral is of polynomial type, by Morales-Ramis theory we can conclude that the connected identity component of the differential Galois group of the first variational equation along any particular solution is an abelian group.

□

8.3. Family III.

THEOREM 8.3. *Consider the family III, The following statements hold.*

- (1) *The first integral of the polynomial vector field is*

$$I = I(x, y) = y - ax^2$$

- (2) *The integral curve of the polynomial vector field is $(x(t), \dot{x}(t))$, where*

$$x(t) = \sqrt{\frac{k_1}{a}} \tan(\sqrt{k_1 a}(k_2 + t)).$$

- (3) *The Differential Galois Group associated to the foliation is isomorphic to the identity group.*

- (4) *The connected identity component of the Differential Galois Group of the first variational equation around any particular solution is an abelian group.*

PROOF. We proceed according to each item,

- (1) The total derivative of $I(x, y)$ vanishes, i.e., $\dot{I} = 0$, therefore I is a first integral of the vector field related to family III.
 (2) Due to $y = \dot{x}$, we obtain $\ddot{x} = a\dot{z}$, where $z = x^2$. Thus, $\dot{x} = ax^2 + k_1$, which implies that

$$\int \frac{dx}{ax^2 + k_1} = t + k_2$$

and then $x(t) = \sqrt{\frac{k_1}{a}} \tan(\sqrt{ak_1}(k_2 + t))$.

- (3) The foliation associated to the vector field of Family II is

$$y' = 2ax, \quad ' := \frac{d}{dx}.$$

Then the solution of this foliation is

$$y(x) = ax^2 + k_1$$

. Then we can conclude that $DGal(L/K)$ has one element, i.e., $DGal(L/K) = I_2$.

- (4) Let $(x_0(t), \dot{x}_0(t))$ be a particular solution of the polynomial vector field related with Family III. Thus, the first variational equation is

$$\frac{d}{dt} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -2ay_0(t) & 2ax_0(t) \end{pmatrix} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix},$$

which is equivalent to

$$\ddot{\xi} - 2ax_0(t)\dot{\xi} - 2ay_0(t)\xi = 0, \quad \xi = \xi_1.$$

Due to the first integral is of polynomial type, by Morales-Ramis theory we can conclude that the connected identity component of the differential Galois group of the first variational equation along any particular solution is an abelian group. □

8.4. Family IV.

THEOREM 8.4. *Consider the family IV, being $p = -4$. The following statements hold.*

- (1) *The dynamical system is hamiltonian with one degree of freedom and with polynomial first integral*

$$H = H(x, y) = \frac{y^2}{2} + \frac{c}{3}x^3 + \frac{3}{4}a^2x^2$$

- (2) *The integral curve of the Hamiltonian vector field is given in terms of P-function.*
 (3) *The Differential Galois Group associated to the foliation is isomorphic to \mathbb{Z}_2 .*
 (4) *The connected identity component of the Differential Galois Group of the first variational equation along any particular solution is an abelian group.*

PROOF. We proceed according to each item,

- (1) The polynomial vector field related with family IV is equivalent to Equation (2.8) being $f(x) = -cx^2 - \frac{3}{2}a^2x$. In virtue of Equation (2.9) we have the Hamiltonian $H = \frac{y^2}{2} + \frac{c}{3}x^3 + \frac{3}{4}a^2x^2$.
- (2) Due to $y = \dot{x}$, we obtain $y^2 = -\frac{2c^2}{3} - \frac{3}{2}a^2x + 2H$. Because previous expression is a cubic polynomial in x , we can do a suitable change of variable to arrive to the elliptic curve given in Equation (2.10) with invariants g_2 and g_3 . Thus, the integral curve of the Hamiltonian system is written in terms of P-function.
- (3) The foliation associated to the vector field of Family IV is

$$y' = -\frac{cx^2 - \frac{3}{2}a^2x}{y}, \quad ' := \frac{d}{dx}.$$

Setting $z = \frac{y^2}{2}$, we obtain $z' = -cx^2 - \frac{3}{2}a^2x$ and therefore $z = -\frac{c}{3}x^3 - \frac{3}{4}a^2x^2$. Due to the differential field K is the field of rational functions $\mathbb{C}(x)$, $\sigma(z) = z$ and $\sigma(y) = \lambda\sqrt{z}$, where $\lambda^2 = 1$. Thus, the Picard-Vessiot extension L is a quadratic extension of K and we can conclude that $DGal(L/K)$ has two elements.

- (4) Let $(x_0(t), \dot{x}_0(t))$ be a particular solution of the polynomial vector field related with Family IV. Thus, the first variational equation is

$$\frac{d}{dt} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -\frac{3}{2}a^2 - 2cx_0(t) & 0 \end{pmatrix} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix},$$

which is equivalent to $\ddot{\xi} = (-\frac{3}{2}a^2 - 2cx_0(t))\xi$, being $\xi = \xi_1$. By Morales-Ramis theory, due to the dynamical system is polynomially integrable, the differential Galois group of this first variational equation is abelian. □

8.5. Family V.

THEOREM 8.5. *Consider the family V, being $s = -4$. The following statements hold.*

- (1) *The dynamical system is hamiltonian with one degree of freedom and with polynomial first integral*

$$H = H(x, y) = \frac{y^2}{2} + \frac{c}{3}x^3 + \frac{3}{4}bx^2$$

- (2) *The integral curve of the Hamiltonian vector field is given in terms of P-function.*
- (3) *The Differential Galois Group associated to the foliation is isomorphic to \mathbb{Z}_2 .*
- (4) *The connected identity component of the Differential Galois Group of the first variational equation along any particular solution is an abelian group.*

PROOF. We proceed according to each item,

- (1) The polynomial vector field related with family V is equivalent to Equation (2.8) being $f(x) = -cx^2 - \frac{3}{2}bx$. In virtue of Equation (2.9) we have the Hamiltonian $H = \frac{y^2}{2} + \frac{c}{3}x^3 + \frac{3}{4}bx^2$.

- (2) Due to $y = \dot{x}$, we obtain $y^2 = -\frac{2c}{3}x^2 - \frac{3}{2}bx + 2H$. Because previous expression is a cubic polynomial in x , we can do a suitable change of variable to arrive to the elliptic curve given in Equation (2.10) with invariants g_2 and g_3 . Thus, the integral curve of the Hamiltonian system is written in terms of P-function.
- (3) The foliation associated to the vector field of Family V is

$$y' = -\frac{cx^2 - \frac{3}{2}bx}{y}, \quad ' := \frac{d}{dx}.$$

Setting $z = \frac{y^2}{2}$, we obtain $z' = -cx^2 - \frac{3}{2}bx$ and therefore $z = -\frac{c}{3}x^3 - \frac{3}{4}bx^2$. Due to the differential field K is the field of rational functions $\mathbb{C}(x)$, $\sigma(z) = z$ and $\sigma(y) = \lambda\sqrt{z}$, where $\lambda^2 = 1$. Thus, the Picard-Vessiot extension L is a quadratic extension of K and we can conclude that $DGal(L/K)$ has two elements.

- (4) Let $(x_0(t), \dot{x}_0(t))$ be a particular solution of the polynomial vector field related with Family V. Thus, the first variational equation is

$$\frac{d}{dt} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -\frac{3}{2}b - 2cx_0(t) & 0 \end{pmatrix} \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix},$$

which is equivalent to $\ddot{\xi} = (-\frac{3}{2}b - 2cx_0(t))\xi$, being $\xi = \xi_1$. By Morales-Ramis theory, due to the dynamical system is polynomially integrable, the differential Galois group of this first variational equation is abelian. \square

9. Conclusion

An in-depth analysis of the quadratic systems containing certain multiparametric families was carried out, for this purpose they were identified and classified, with the aim of making more bearable the study on the stability of its critical points both in the finite and infinite plane. The existence of transcritical bifurcations in the given system was determined. Finally, a study was made on the hamiltonian cases and the differential Galois groups of their foliations and variational equations.

References

- [1] ACOSTA-HUMÁNEZ P.B., LAZARO J.T., MORALES-RUIZ J.J. & PANTAZI CH., *On the integrability of polynomial fields in the plane by means of Picard-Vessiot theory*, Discrete & Continuous Dynamical Systems-A **35** (2015): 1767–1800.. Available at arXiv:1012.4796.
- [2] ACOSTA-HUMÁNEZ P. B., REYES LINERO A. & RODRÍGUEZ CONTRERAS J., *Algebraic and qualitative remarks about the family $yy' = (\alpha x^{m+k-1} + \beta x^{m-k-1})y + \gamma x^{2m-2k-1}$* , preprint 2014. Available at arXiv:1807.03551.
- [3] RODRÍGUEZ CONTRERAS J., ACOSTA-HUMÁNEZ P. B. & REYES LINERO A., *Algebraic and qualitative remarks about the family $yy' = (\alpha x^{m+k-1} + \beta x^{m-k-1})y + \gamma x^{2m-2k-1}$* , Open Mathematics **17** (2019), 1220–1238.
- [4] ACOSTA-HUMÁNEZ P. B., REYES LINERO A. & RODRÍGUEZ CONTRERAS J., *Galoisian and Qualitative Approaches to Linear Polyanin-Zaitsev Vector Fields*, Open Mathematics **16** (2018), 1204–1217. Available at arXiv:1807.05272.
- [5] Acosta-Humánez P. B., Campo Donado M., Reyes Linero A., & Rodríguez Contreras J., (2019). Algebraic and qualitative aspects of quadratic vector fields related with classical orthogonal polynomials. arXiv preprint arXiv:1906.09764.
- [6] Rodríguez Contreras J, Reyes Linero A., Campo Donado M., & Acosta-Humánez P. B. , (2020). Dynamical and Algebraic Analysis of Planar Polynomial Vector Fields Linked to Orthogonal Polynomials. Journal of Southwest Jiaotong University, 55(4).

- [7] TORRES HENAO J. A. , *Sistemas Dinámicos Planos*. Universidad Nacional de Colombia, Facultad de Ciencias, Escuela de Matemáticas, Medellín, Colombia (2013). <http://bdigital.unal.edu.co/9478/>
- [8] VÍLCHEZ LOBATO M. L. , VELASCO MORENTE F., GARCÍA DEL HOYO J. J., *Bifurcaciones transcriticals y ciclos límites en un modelo dinámico de competición entre dos especies. Una aplicación a la pescadera de engraulis encrasicolus de la Región Suratlántica española*, (2002).
- [9] XIANGDONG, X., & JIANFENG, Z., *Plane Polynomial System and it's Accompany System*, Journal of Mathematical Study, 37(2) (2004), 161–166.
- [10] GAIKO V.A., *Multiple limit cycle bifurcations of the FitzHugh–Nagumo neuronal model*, Nonlinear Analysis: Theory, Methods & Applications, 74(18), (2011) 7532–7542
- [11] PERKO, L., *Differential Equations and Dynamical Systems (New York: Springer Verlag)*, 2001.
- [12] GUCKEINHEIMER, J., Y HOLMES, P., *Nonlinear Oscillations, Dynamical Systems, and Bifurcations of Vector Fields (New York: Springer Verlag)*, 1983
- [13] DUMORTIER F., LLIBRE J. & ARTES J., *Qualitative Theory Of Planar Differential Systems*, (Berlin: Springer), 2006.
- [14] HALE J.K & KOÇAK H., *Dynamics and Bifurcations*, (Springer-Verlag), 1991.
- [15] HERSENS C., MAESSCHALCK P., ARTÉS J. C., DUMORTIER F. & LLIBRE J., *P4* (<http://mat.uab.es/artes/p4/p4.htm>), Dept. de Matemàtiques, Universitat Autònoma de Barcelona.
- [16] COMPUTATIONAL ALGEBRAIC SYSTEM, DYNAMIC GEOMETRY SOFTWARE; GEOGEBRA.
- [17] Acosta-Humánez, P. & Jiménez G. *Some tastings in Morales-Ramis theory*, Journal of Physics: Conference Series. Vol. 1414. No. 1. IOP Publishing, 2019.
- [18] Acosta-Humánez, M. F., Acosta-Humánez, P. B., & Tuirán, *Generalized Lennard-Jones Potentials, SUSYQM and Differential Galois Theory*, SIGMA. Symmetry, Integrability and Geometry: Methods and Applications, 14, (2018) 099.
- [19] Acosta-Humánez, P. B., Blázquez-Sanz, D., & Vargas-Contreras, C. A. *On Hamiltonian potentials with quartic polynomial normal variational equations*, Nonlinear Studies, 16(3), (2009) 299–314.
- [20] Acosta-Humánez, P., & Blázquez-Sanz, D. *Non-integrability of some hamiltonians with rational potentials*, Discrete & Continuous Dynamical Systems-B, 10(2&3), (2008) 265–293.
- [21] Acosta-Humánez, P. B. *La teoría de Morales-Ramis y el algoritmo de Kovacic. Lecturas Matemáticas*, 27(3), (2006) 21–56.
- [22] Acosta-Humánez, P. B. *Nonautonomous Hamiltonian Systems and Morales–Ramis Theory I. The case $\ddot{x} = f(x, t)$* , SIAM Journal on Applied Dynamical Systems, 8(1), (2009) 279–297.
- [23] Acosta-Humánez, P. B., & Pantazi, C. *Darboux integrals for Schrödinger planar vector fields via Darboux transformations*, SIGMA. Symmetry, Integrability and Geometry: Methods and Applications, 8 (2012) 043.
- [24] Acosta-Humánez, P. B., Álvarez-Ramírez, M., & Delgado, J. *Non-integrability of some few body problems in two degrees of freedom*, Qualitative theory of dynamical systems, 8(2), (2009) 209–239.
- [25] Acosta-Humánez, P. B., Alvarez-Ramírez, M., Blázquez-Sanz, D., & Delgado, J. *Non-integrability criterium for normal variational equations around an integrable subsystem and an example: The Wilberforce spring-pendulum*, Discrete & Continuous Dynamical Systems-A, 33(3), (2013) 965–986.
- [26] Acosta-Humánez, P. B., Lázaro, J. T., Morales-Ruiz, J. J., & Pantazi, C. *Differential Galois theory and non-integrability of planar polynomial vector fields*, Journal of Differential Equations, 264(12), 7183–7212.
- [27] Acosta-Humánez, P. B., & Blázquez-Sanz, D. A. V. I. D. *Hamiltonian system and variational equations with polynomial coefficients* Dynamic systems and applications, Dynamic, Atlanta, GA, 5, (2008) 6–10.
- [28] Acosta-Humánez, P. B., Blázquez-Sanz, D., & Venegas-Gómez, H. *Liouvillian solutions for second order linear differential equations with polynomial coefficients*, São Paulo Journal of Mathematical Sciences, (2020) 1–20.
- [29] Acosta-Humánez, P. B. *Métodos algebraicos en sistemas dinámicos*, Ediciones Universidad del Atlántico, EMALCA 2014.
- [30] Acosta-Humánez, P. B., & Yagasaki, K. *Nonintegrability of the unfoldings of codimension-two bifurcations*, Nonlinearity, 33(4), (2020) 1366–1387.

[31] Abramowitz, M. and Stegun, I. Handbook of Mathematical Functions, ninth printing, New York: Dover (1972).

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